

Spectroscopic constrains on variation of fundamental constants in astrophysics

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New ideas in low-energy tests of fundamental physics

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Sesto Conference

*Varying fundamental constants and
dynamical dark energy*

July 2013 (Italy)



Fundamental constants in atomic physics

Fundamental constants, which influence atomic and molecular spectra:

- Fine structure constant $\alpha = e^2/(\hbar c)$ is a coupling constant in QED.
- Electron to proton mass ratio $\mu = m_e/m_p$. Because m_p is proportional to Λ_{QCD} , $\mu \sim m_e/\Lambda_{QCD}$.
- Nuclear gyromagnetic ratio g_n can be expressed in terms of Λ_{QCD} and quark masses, but for atomic physics g_n is independent constant. It **always** enters in combination $g_n\mu$. According to Flambaum & Tedesco (2006) the dependence of g_n on quark masses is **weak**.

Dimensionless sensitivity coefficients

If fundamental constants change, the frequency of any atomic transition also change:

$$\begin{aligned}\omega &= \omega_0 \left[1 + Q_\alpha \frac{\delta\alpha}{\alpha} + Q_\mu \frac{\delta\mu}{\mu} + Q_g \frac{\delta g_n}{g_n} \right], \\ \frac{\delta\omega}{\omega} &= \frac{\delta F}{F}, \quad F = \alpha^{Q_\alpha} \mu^{Q_\mu} g_n^{Q_g}.\end{aligned}$$

In order to detect this variation we need to compare at least two transition frequencies:

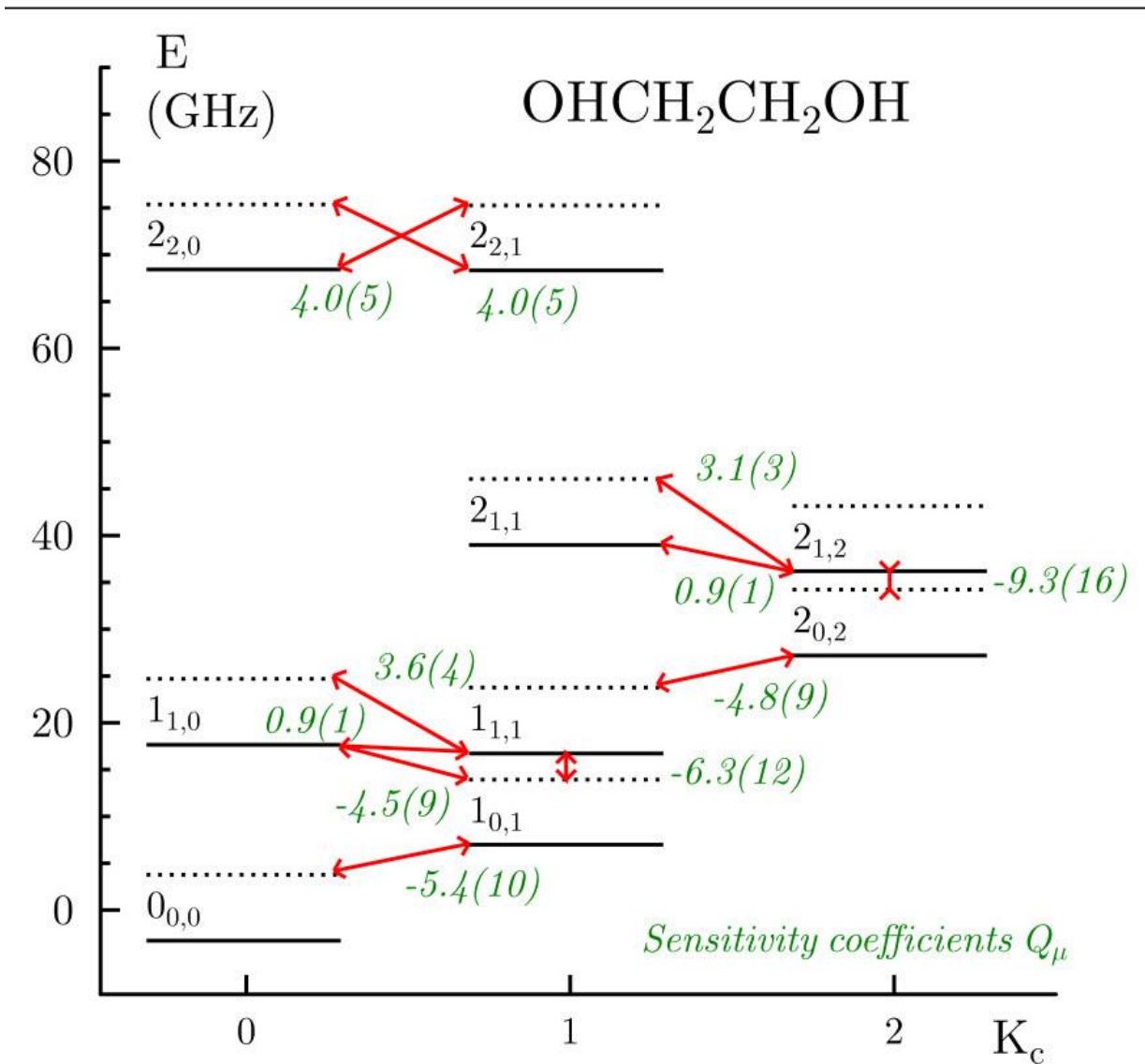
$$\frac{\omega_i}{\omega_k} = \left(\frac{\omega_i}{\omega_k} \right)_0 \left[1 + \Delta Q_\alpha \frac{\delta\alpha}{\alpha} + \Delta Q_\mu \frac{\delta\mu}{\mu} + \Delta Q_g \frac{\delta g_n}{g_n} \right].$$

Sensitivity coefficients for different wavebands (in a.u.)

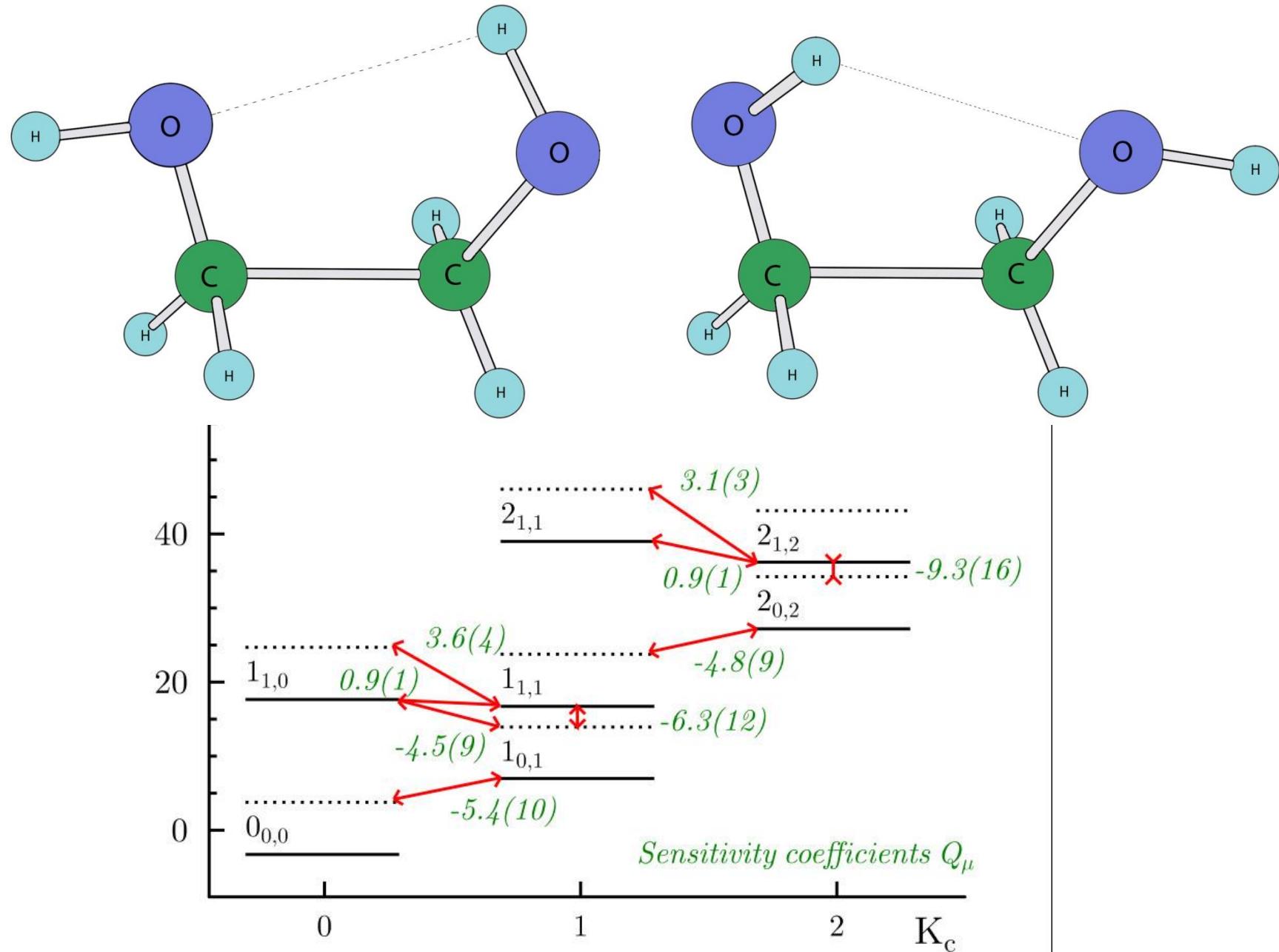
- For optical transitions in atoms and molecules with nuclear charge $Z \leq 30$, all sensitivities are small, $Q_\alpha, Q_\mu, Q_g \ll 1$.
- Optical transitions in Highly Charged Ions: $|Q_\alpha| \gg 1$.
- Fine structure (IR, FIR): $\sim \alpha^2 \Rightarrow Q_\alpha = 2$.
- Vibrational structure (IR): $\sim \mu^{1/2} \Rightarrow Q_\mu = \frac{1}{2}$.
- Rotational structure (FIR, microwave): $Q_\mu = 1$.
- Magnetic hyperfine structure (microwave):
 $Q_\alpha = 2$; $Q_\mu = 1$; $Q_g = 1$.
- Tunneling transitions in polyatomic molecules (FIR, microwave):
 $1 \lesssim Q_\mu \lesssim 10$.
- Microwave mixed tunneling-rotational lines: $|Q_\mu| \gg 1$.
- Microwave Λ -doublet, Ω -doublet, and K -doublet lines in linear radicals: $|Q_\alpha|, |Q_\mu| \gg 1$.

Tunneling-rotational spectrum of Ethylene glycol

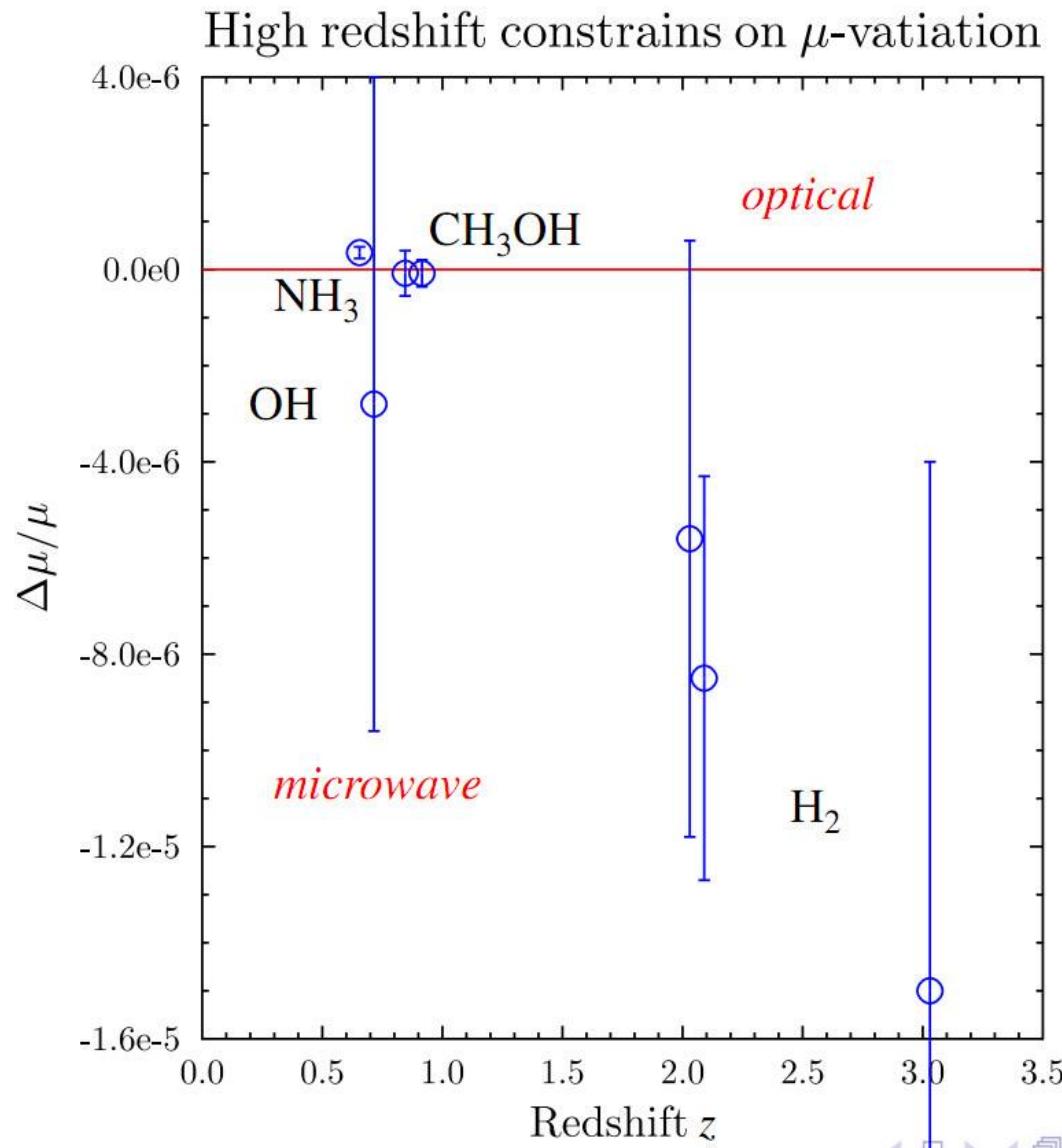
[A Viatkina & MK 2014]



Tunneling-rotational spectrum of Ethylene glycol

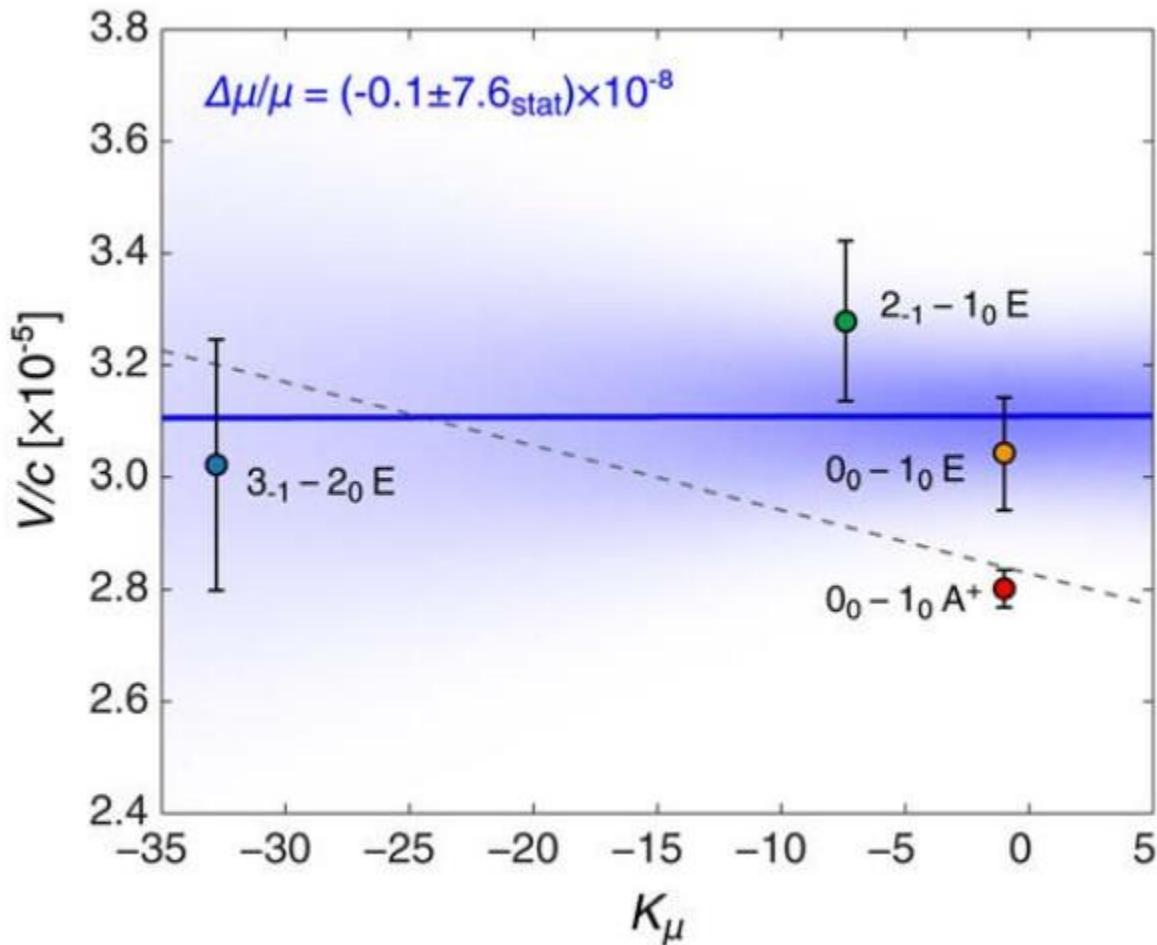


Ограничения на вариацию μ на больших z



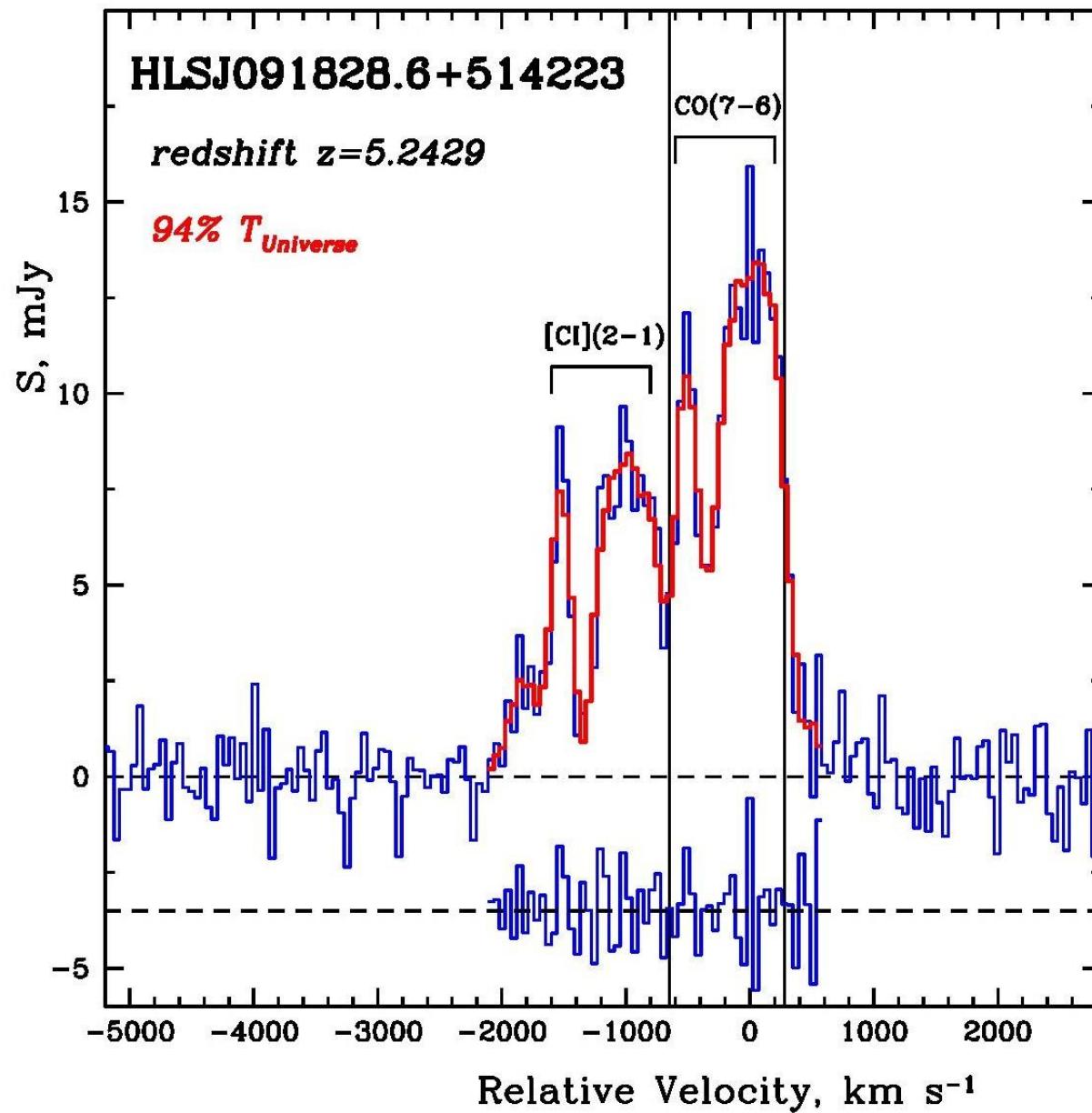
Constrain on μ -variation from observation of methanol lines at redshift $z=0.89$

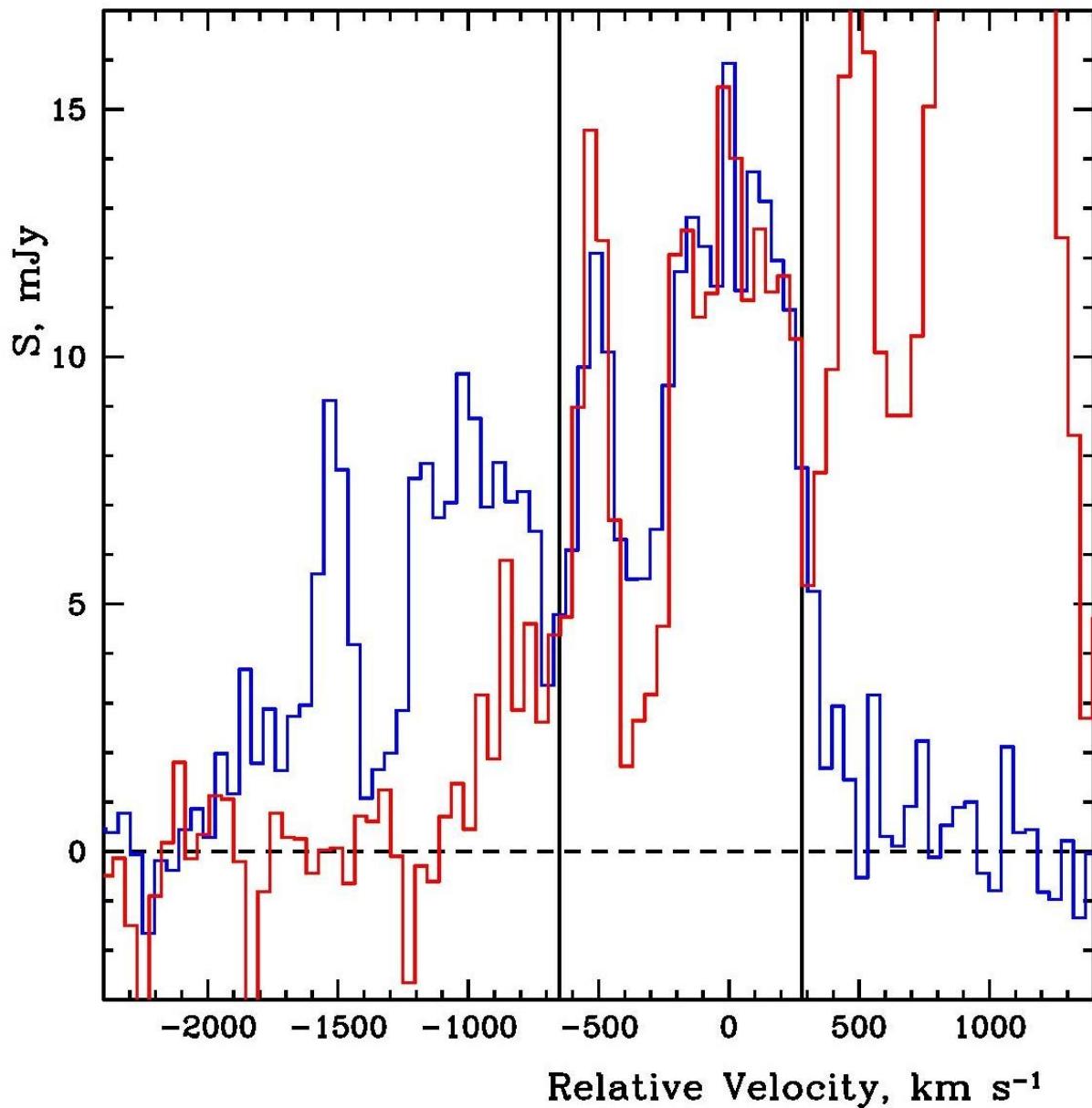
[Bagdonaitė et al. Science 339, 46 (2013)]

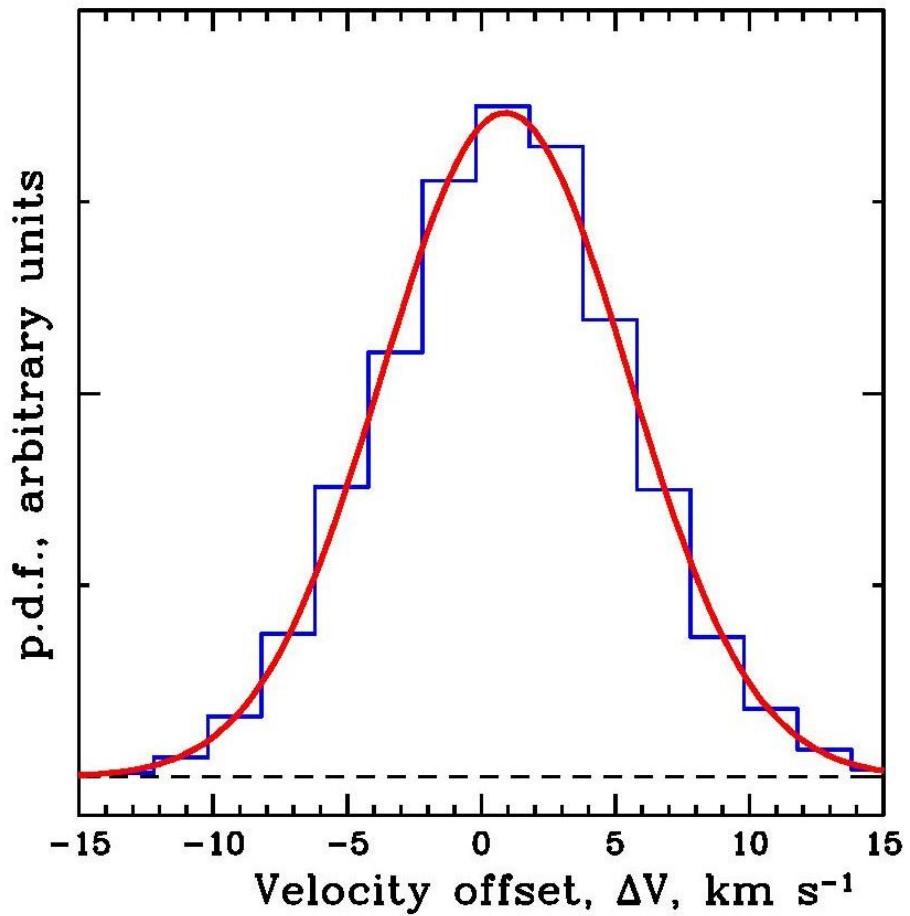


3P_2 - 3P_1 fine structure transition in C I
at the redshift z=5.2

(S. A. Levshakov et al *Astron. Astrophys.*,
2012, 540, L9)







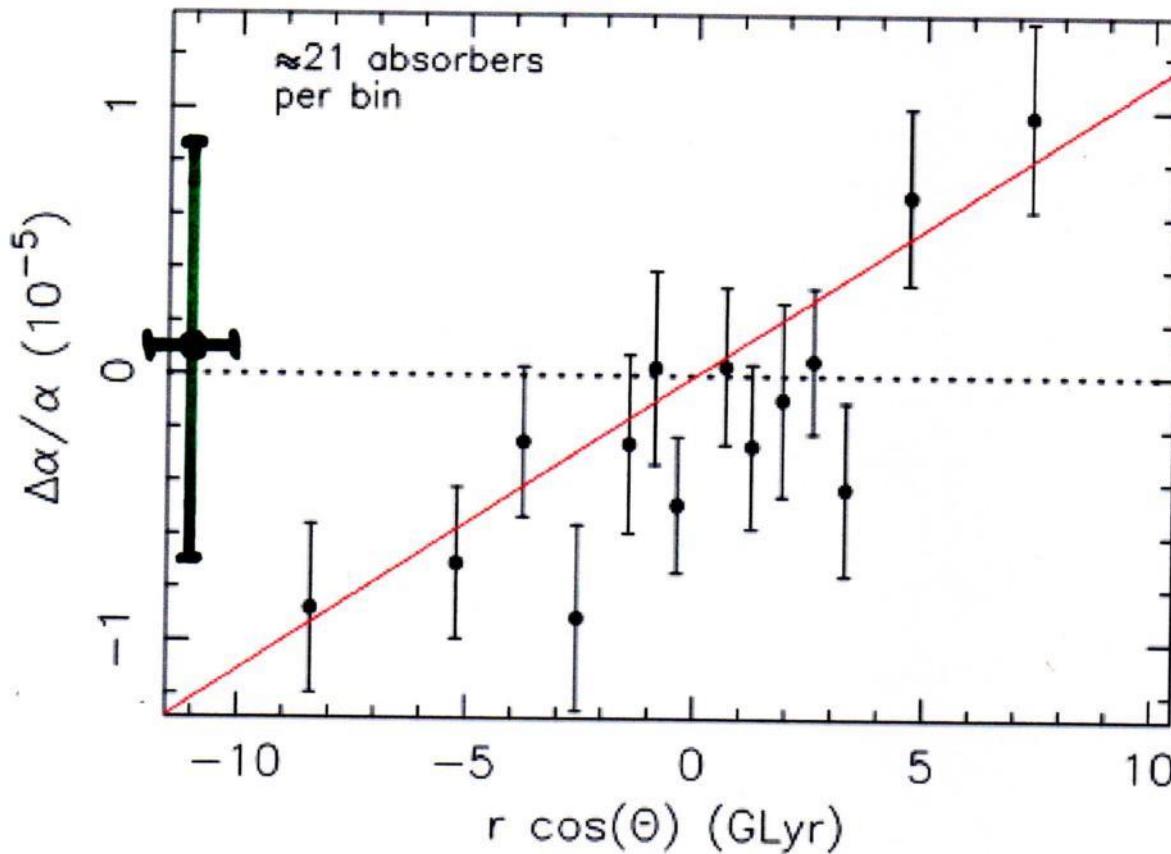
$$\Delta V = (1 \pm 5) \text{ km/s}$$

$$F = a^2 / \mu$$

$$\Delta F/F < 2 \times 10^{-5}$$

Australian dipole

(Webb et al, *Phys. Rev. Lett.*, 2011, 107, 191101)



NH_3 and CH molecules in cold interstellar clouds in the Milky Way

*Testing chameleon models of the
Dark energy*

Tunneling transition in NH_3 is highly sensitive to μ -variation, while Ω -doublet transitions in CH are sensitive to variation of both μ and α .

Radio astronomical observations of NH₃

[Levshakov et al, 2014]

32m MEDICINA
Italy



NH₃, HC₃N

100m EFFELSBERG
Germany



NH₃, HC₃N,
HC₅N, HC₇N

45m NOBEYAMA
Japan

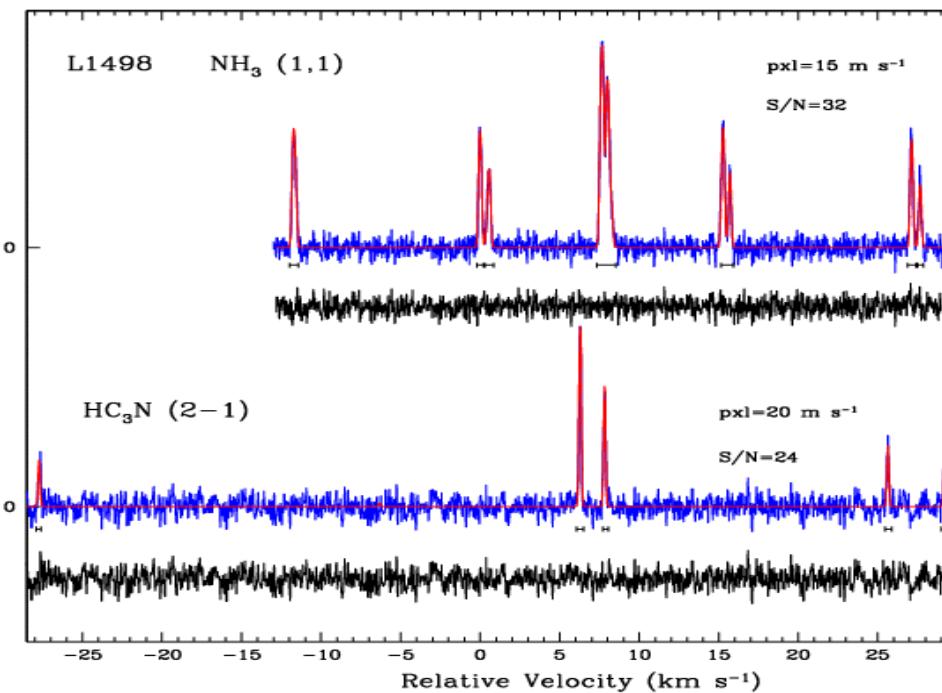


NH₃, N₂H⁺

41 molecular cores
in Taurus

52 molecular cores
in Aquila

Effelsberg 100-m



Line width:

$\text{FWHM}=200 \text{ m/s}$

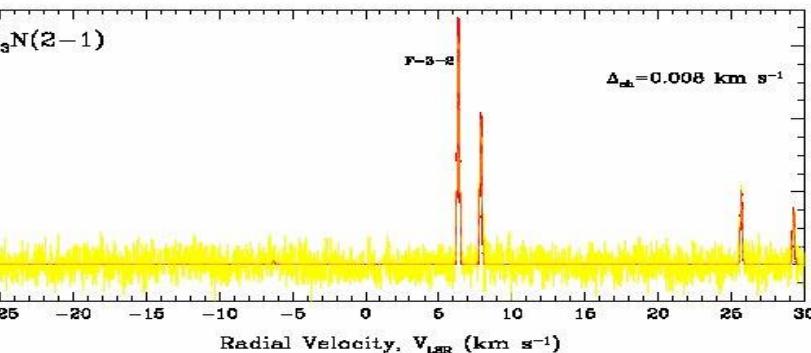
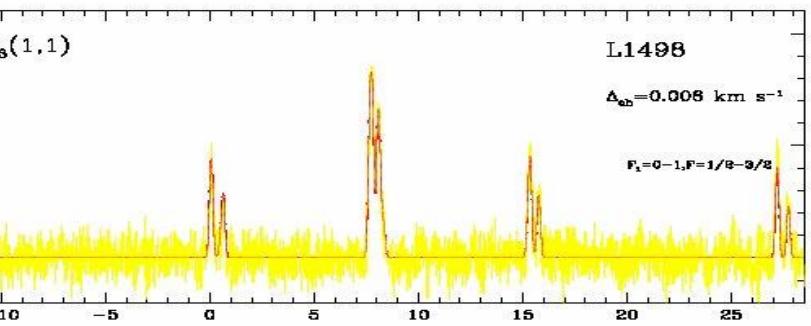
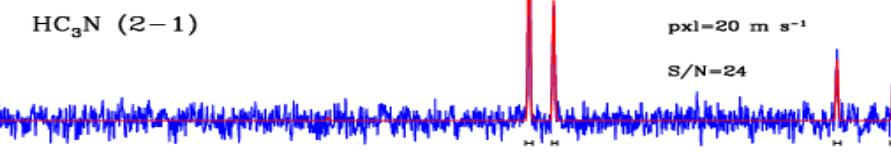
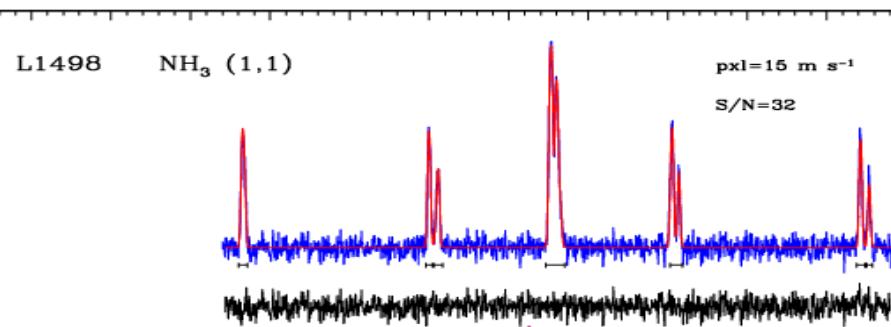
$$\sigma \sim 1 \text{ m/s}$$

Line position
uncertainty

$\text{FWHM}=150 \text{ m/s}$

$$\sigma \sim 5 \text{ m/s}$$

Effelsberg 100-m



Line width:

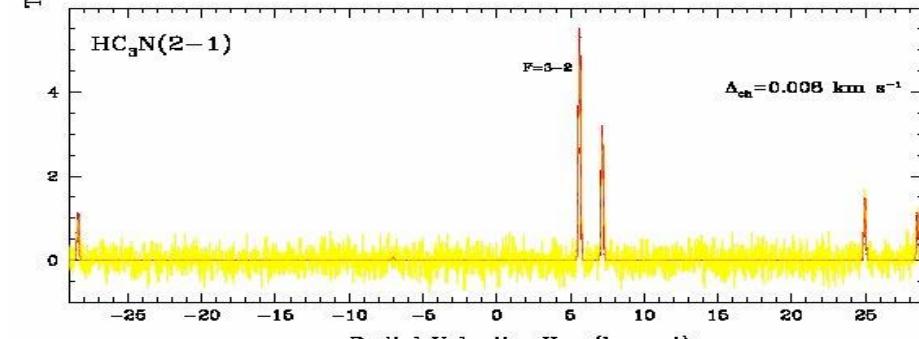
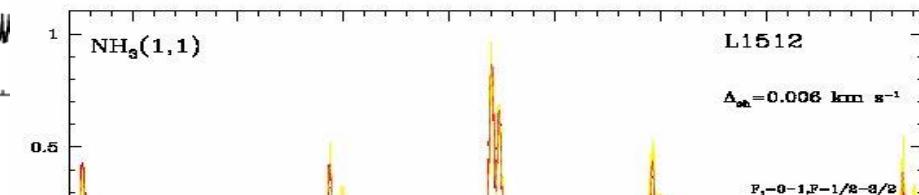
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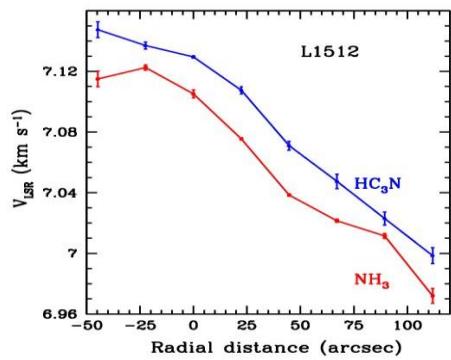
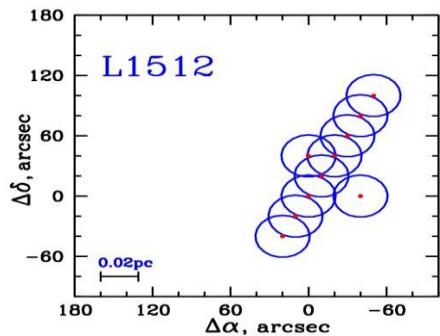
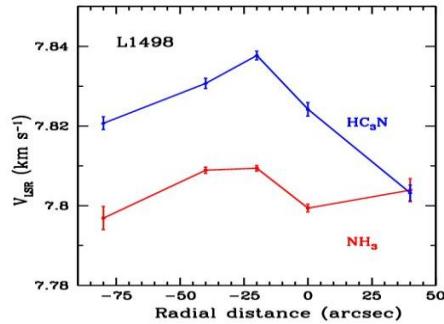
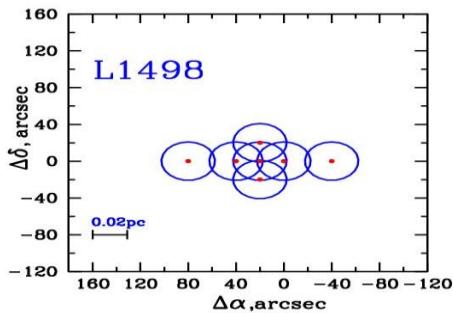
Line position uncertainty

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Effelsberg, 2010 mapping



L1498+L1512

Jan, 2010:

$$\Delta V = 27 \pm 1 \pm 3 \text{ m/s}$$

$$\Delta \mu/\mu = 26 \pm 1 \pm 3 \text{ ppb}$$

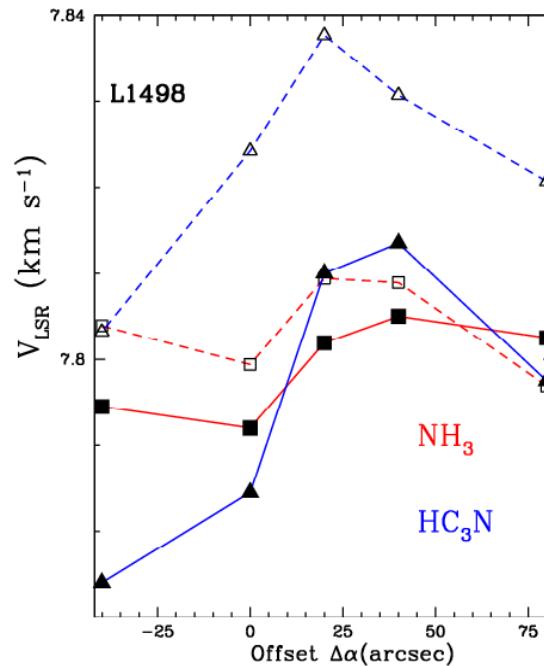
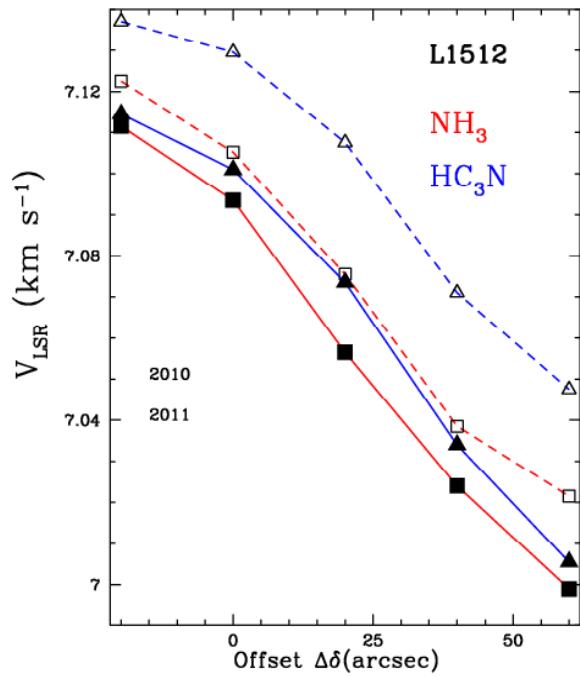
Feb, 2009:

$$\Delta V = 26 \pm 4 \pm 3 \text{ m/s}$$

$$\Delta \mu/\mu = 26 \pm 4 \pm 3 \text{ ppb}$$

additional tests: data reproducibility

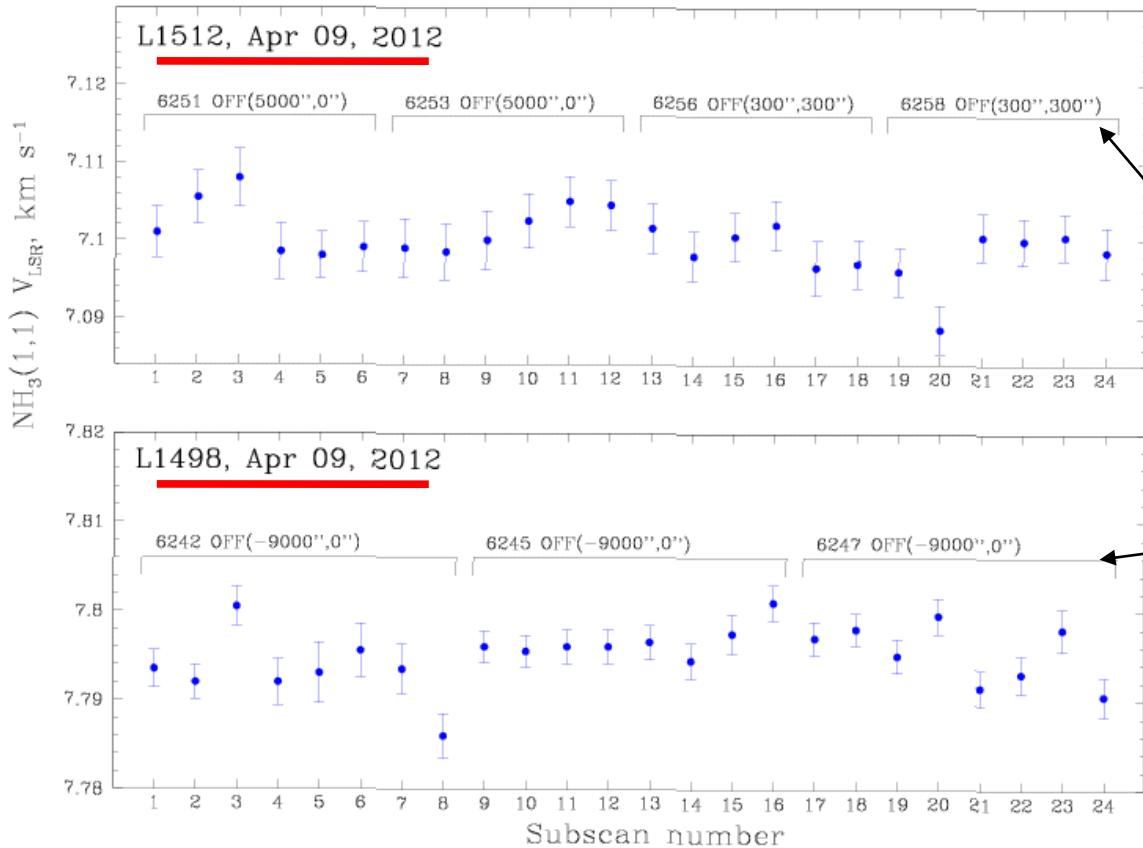
Effelsberg, 2011
mapping



FFTS (fast fourier transform spectrometer)

Time series

Effelsberg, 2012



New spectrometer:

XFFTS
(eXtende
d FFTS)

Exposure
time:

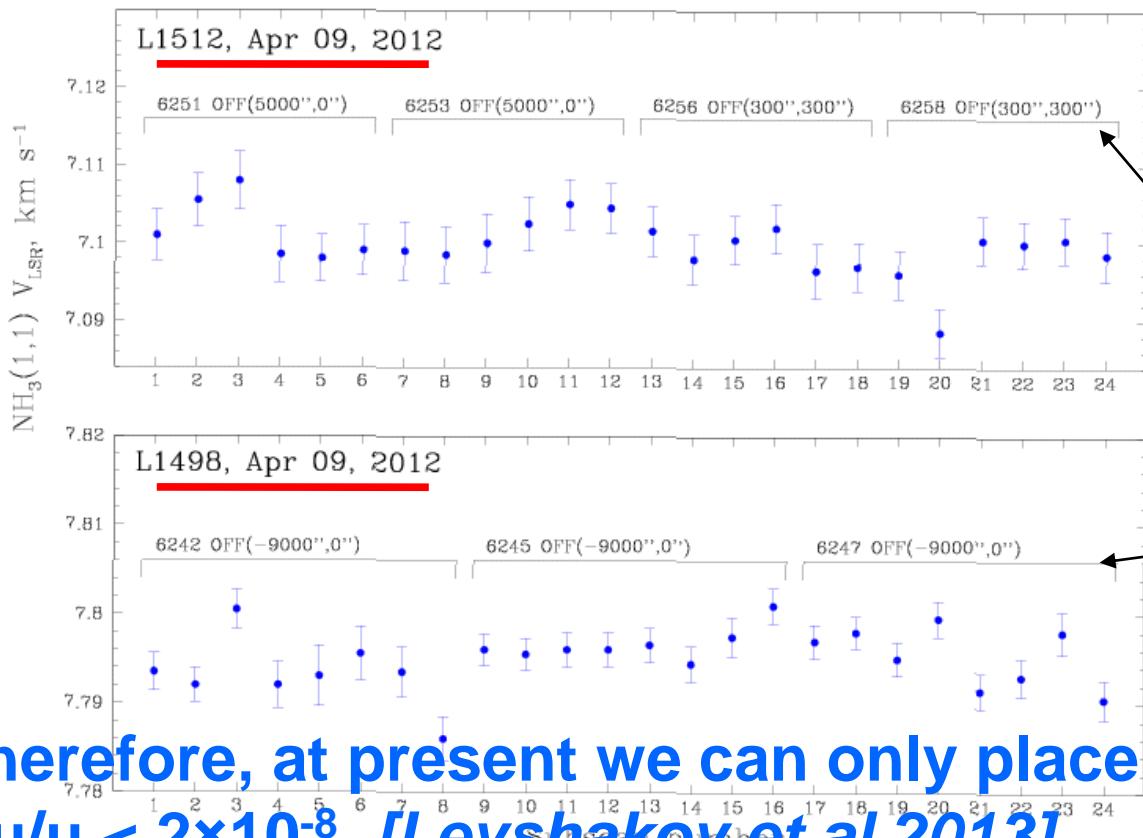
30 min/scan
(ON+OFF)
PSW
150 sec/point

instability of $\delta V \sim 10$ m/s detected

independently checked by Benjamin Winkel (MPIfR)

Time series

Effelsberg, 2012



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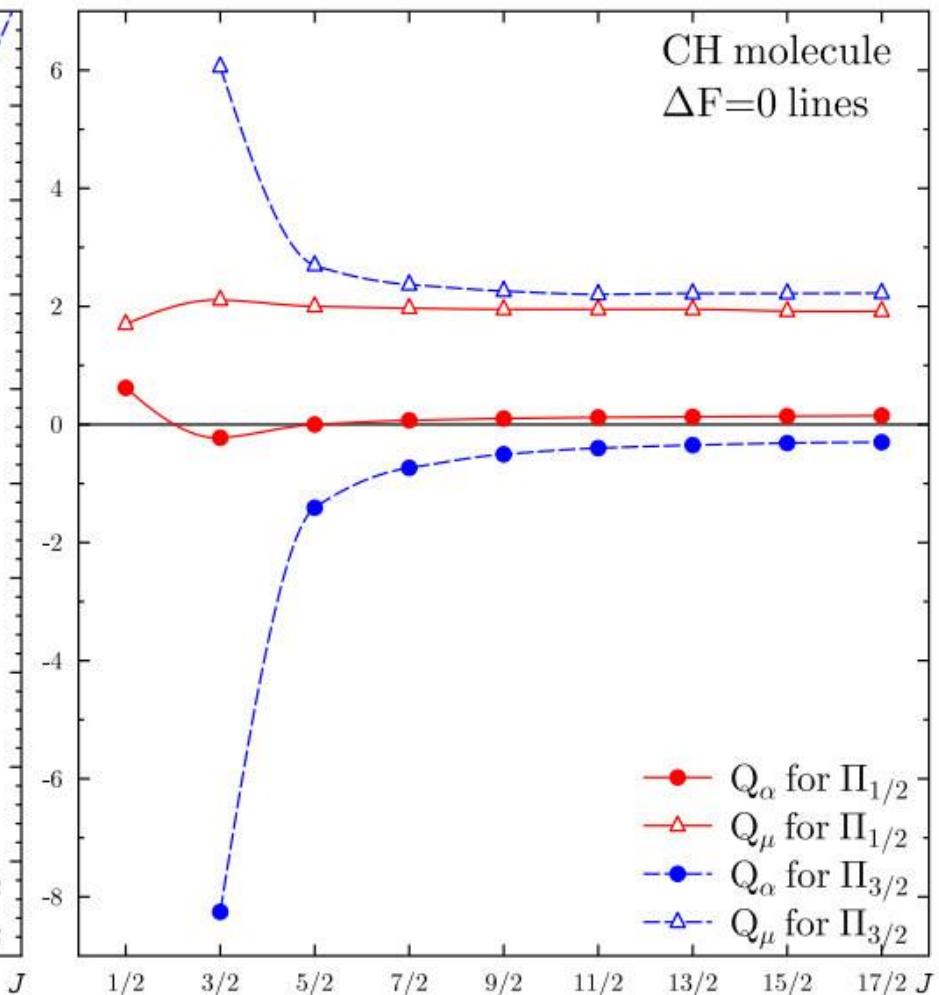
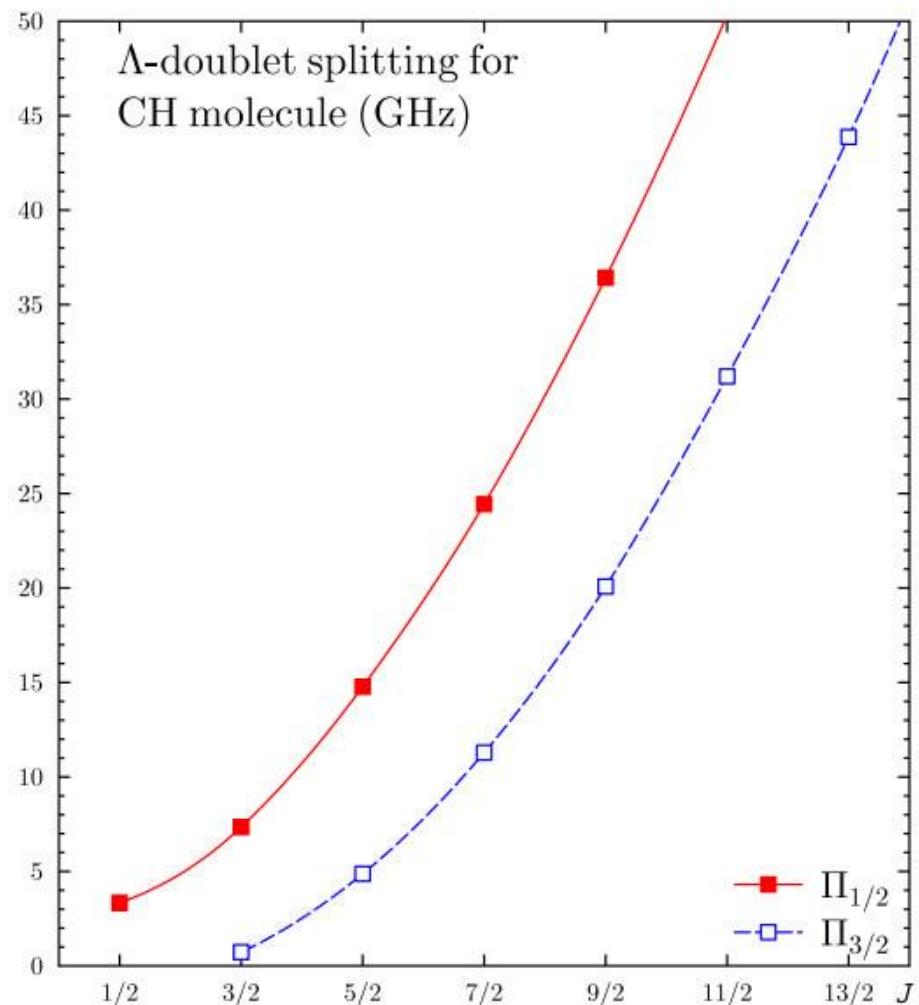
Therefore, at present we can only place an upper bound:
 $\Delta\mu/\mu < 2 \times 10^{-8}$ [Levshakov et al 2013]

instability of $\delta V \sim 10 \text{ m/s}$ detected

independently checked by Benjamin Winkel (MPIfR)

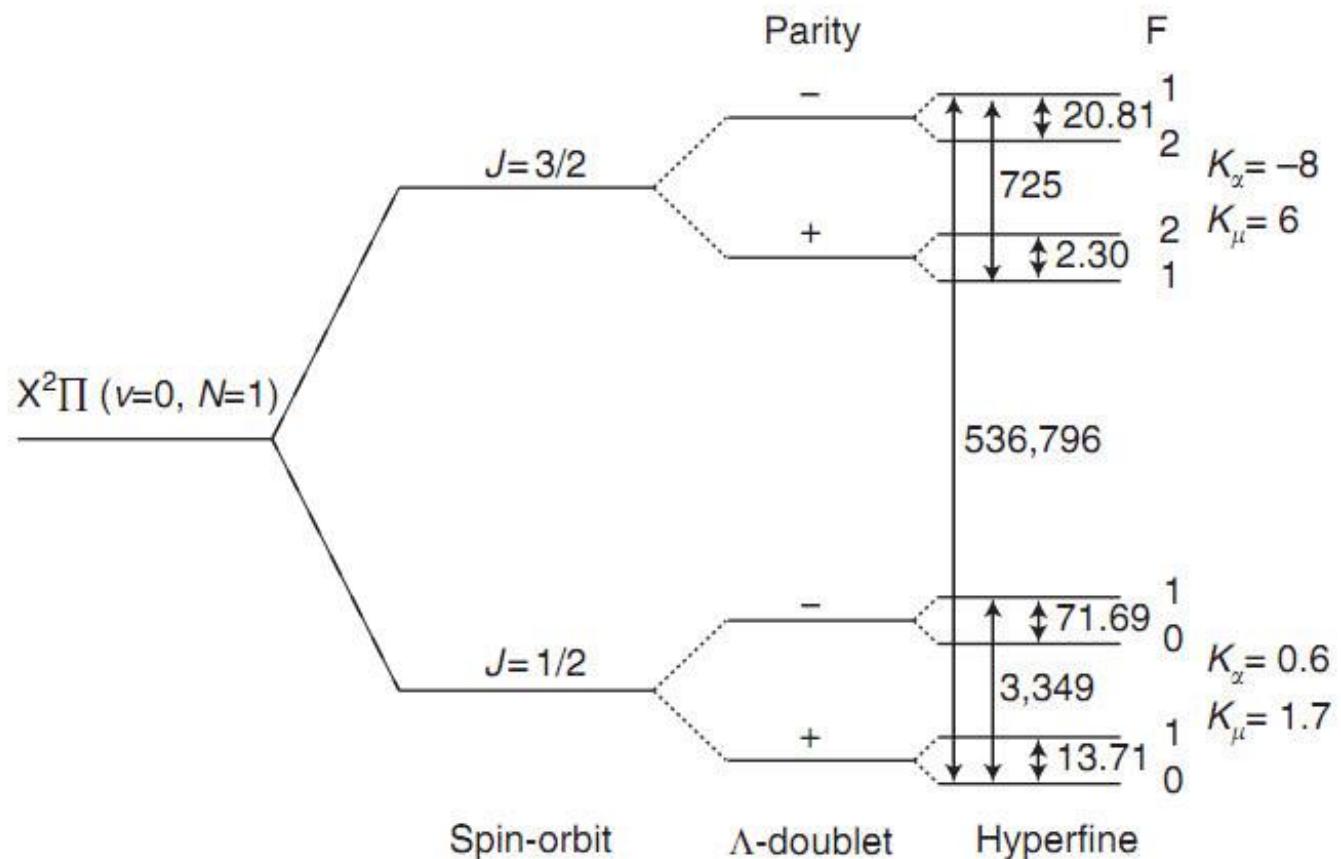
Λ -doublet, or Ω -doublet transitions in CH

Spin-orbit interaction couples electron spin to the molecular axis. When rotational energy grows, electron spin decouples from the axis. Then quantum number Ω is substituted by Λ . Competition between Coriolis and SO interactions leads to strong dependence of the doubling splitting on α and μ .

Frequencies & sensitivities of Λ -transitions in CH

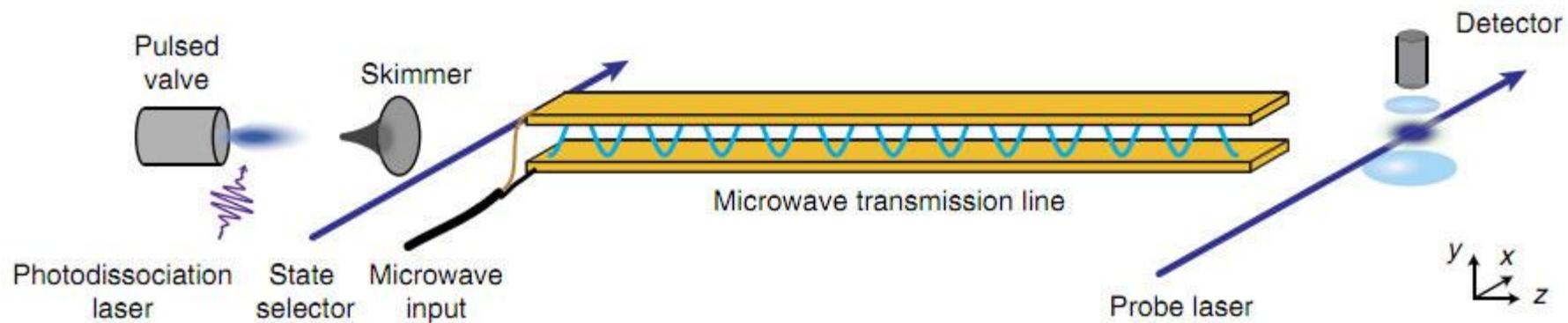
Detection of Λ -doublet transitions in CH

[Truppe et al, Nat. Commun. 4, 2600 (2013)]



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Table 1 | Measured Λ -doublet frequencies with 1σ uncertainties.

Transition	Frequency (Hz)
$(1/2^+, 1) - (1/2^-, 1)$	$3,335,479,356 \pm 3$
$(1/2^+, 0) - (1/2^-, 1)$	$3,349,192,556 \pm 3$
$(1/2^+, 1) - (1/2^-, 0)$	$3,263,793,447 \pm 3$
$(3/2^+, 2) - (3/2^-, 2)$	$701,677,682 \pm 6$
$(3/2^+, 1) - (3/2^-, 1)$	$724,788,315 \pm 16$
$(3/2^+, 1) - (3/2^-, 2)$	$703,978,340 \pm 21$
$(3/2^+, 2) - (3/2^-, 1)$	$722,487,624 \pm 16$

Levels are labelled with the notation (J^P, F) .

Table 2 | Analysis of astronomical data.

Source	Transition 1	Transition 2
G111.7 – 2.1(CasA)	CH(3264, 3335, 3349)	OH(1667)
G265.1 + 1.5(RCW36)	CH(3264, 3335)	OH(1612, 1665, 1667, 1721)
G174.3 – 13.4(Heiles2)	CH(3264, 3335, 3349)	OH(1612, 1665, 1667, 1721)
G6.0 + 36.7(L134N)	CH(3264, 3335, 3349)	OH(1665, 1667)
G49.5 – 0.4(W51)	CH(702)	CH(3264, 3335, 3349)

v (km s $^{-1}$)	Δv_{12} (km s $^{-1}$)	$\Delta v'_{12}$ (km s $^{-1}$)	$\frac{\Delta z}{z} (10^{-7})$	$\frac{\Delta \mu}{\mu} (10^{-7})$	Ref.
–1.4, 0	–0.01 (0.09)	–0.08 (0.11)	1.5 (2.0)	–3.1 (4.1)	30,32
6.8	0.06 (0.19)	0.04 (0.16)	0.9 (3.1)	1.9 (6.4)	34
5.8	0.00 (0.19)	–0.02 (0.19)	0.6 (3.6)	–1.2 (7.4)	32,43
2.5	0.05 (0.13)	–0.12 (0.13)	2.3 (2.4)	–4.8 (5.0)	32,43
65	–0.85 (0.53)	–0.48 (0.55)	–1.8 (2.0)	3.6 (4.1)	36

Conclusions

- In the recent years there was gradual shift of emphasis from optical to microwave waveband in the quest for the variation of the fundamental constants.
- At present there is no reliable evidence of the variation of constants either in space, or in time.
- Astrophysical data leads to very strict upper limits on the variation of constants.
- State of the art laboratory techniques are essential to provide high accuracy rest frame frequencies of important molecular transitions.
- New observations of microwave spectra at high redshifts are likely to lead to even stronger limits in the near future.